# **Photodiode Front Ends** *The* **REAL** *Story*

# **SAINT E NICKOLS III**

**Photodiode front ends are by no means glamorous. Living between the detector and the signal processing system, they're supposed to turn a photocurrent into a buffered, filtered electronic replica while preserving the signal-to-noise ratio (SNR). Nobody notices them until they stop doing their jobs. Your optical system may be a thing of great beauty, but a badly designed front end can sink those precious photoelectrons deep in Johnson noise. My unscientific sampling suggests that unfortunately a great many front ends are badly designed; the usual mistake is to trade SNR for speed without a fight. This article will describe techniques for building fast front ends without sacrificing SNR.**

**The State of State Shot-noise level was published.<sup>1</sup> In the** article I wrote on how to get rid of laser intensity noise down to the shot-noise level was published.<sup>1</sup> In the intervening years, in my work with ambient-light systems, I've identified a similar need. There's so much misinformation circulating about photodiode front ends, especially transimpedance amplifiers and their "inherent" high frequency noise peak, that it's time to put the record straight: it really is possible to do fast measurements, *at* the shot noise limit, *at* low light intensities, *with* ordinary components.2 To avoid confusing units, in this article I quote signal-to-noise ratios in terms of electrical power, in a 1-Hz bandwidth, so that power and power spectral density are numerically equal. Also, I'm going to use SNR rhetorically even though I'm computing it by dividing the total signal power by the 1-Hz noise, which is really a carrier-to-noise ratio (CNR).

# **The simplest front end: a resistor**

Let's say we're building an instrument that needs a 1-MHz bandwidth, with a Si PIN photodiode of 100 pF  $C_d$  producing a 2-µA photocurrent  $i_d$ , and that we want to stay shot-noise limited because the background light is quiet. Our first thought is to turn the photocurrent into a voltage

with a resistor. To make the signal swing conveniently large, we might pick 1  $M\Omega$ , as shown in Fig. 1. This circuit is linear, but extremely slow—its 3 dB corner  $f_c =$  $1/(2\pi R_{L}C_{d}) \approx 1600$  Hz, a factor of 600 slower than our design point (applying reverse bias reduces  $C_d$  by as much as 7:1—otherwise it might be 250 Hz). Although the signal rolls off at 1600 Hz, surprisingly enough the signal-to-noise ratio does not deteriorate at all, remaining constant at  $i_d^{\,2}/(i_{N\text{th}}^{\,2}+i_{N\text{s}}^{\,2})$ . The resistor's Johnson noise current  $i_N$  and the photocurrent shot noise  $i_{Ns}$  are both treated exactly as the signal is. The reason is that the signal and noise sources are all in parallel. Thus they all roll off together with frequency, and their ratios are constant, as Fig. 2 shows.

# **Reducing the load resistance**

Reducing  $R_t$  will reduce the  $RC$  product and speed things up. Unlike the *RC* rolloff, this *does* reduce the SNR. The noise current of  $R_i$  is

$$
i_N = \sqrt{\frac{4kT}{R_L}}\tag{1}
$$

so it goes up as  $R<sub>L</sub>$  goes down. Still, we can safely reduce  $R_t$  as long as shot noise dominates. The shot noise of a photocurrent  $I_d$ is

$$
i_N = \sqrt{2 e i_d}
$$
 (2)  
Short noise ceases to dominate when these

become equal, i.e. when  $i_dR_l = 2kT/e$  (51) mV at 300 K), but the SNR loss is less than 1 dB when  $i_aR_i \geq 200$  mV. If that 1 dB is acceptable, we'll choose  $R_l = 200$  mV/2  $\mu$ A = 100 k $\Omega$ , raising  $f_c$  to 16 kHz. This is still much too slow, because the full signal swing appears across  $C_{\phi}$  which hogs all the signal current. Eliminating the swing eliminates the capacitive current, but requires a low-impedance load. How can we avoid degrading the noise?

## **The transimpedance amplifier**

The usual way is to connect the photodiode to virtual ground, as shown in Fig. 3. Although the inverting input of  $A_1$  draws no current, feedback forces the voltage there to be close to zero at all times. The way this works is that  $A<sub>i</sub>$  senses the voltage across  $C_d$ , and wiggles the other end of  $R_f$ to zero it out. Provided  $A<sub>1</sub>$  has high open loop gain  $A_{VOL}$ , the swing across  $C_d$  is greatly reduced, and the bandwidth greatly improved. The amplifier's own capacitance  $C_{\text{in}}$  (2-20 pF) must be added to  $C_d$ . This circuit has been analyzed over and over again in the literature, so we'll just exhibit the results, but in practice we have to pay more attention to frequency compensation.<sup>3</sup> A very good operational amplifier for low speed transimpedance amps is the LF356 (which is unfortunately becoming harder to get). Because the *RC* and the op



**Figure 1**.The world's simplest front end: a load resistor.



**Figure 2**. Photodiode/load resistor circuit: frequency response and 1 Hz SNR.



**Figure 3**.Transimpedance amplifier schematic and noise model.



**Figure 4.** Noise performance of the transimpedance amplifier of Fig. 3, showing the dominance of  $e_{NAmp}$  at high frequency. $A_1$  is an LF356, R<sub>f</sub>=100k $\Omega$ , C<sub>f</sub>=0.5 pF.

amp gain both roll off as 1/*f*, and the loop gain goes as their product, the unity gain crossover of the transimpedance amp moves to about

$$
f_{CL} \approx \sqrt{f_{RC} f_T} \tag{3}
$$

which for the LF356/100 k $\Omega$ /100 pF combination is (16 kHz • 4 MHz)<sup> $1/2 \approx 250$  kHz.</sup> The transimpedance rolls off somewhat earlier than this, since it depends on the magnitudes of the impedances of the feedback elements, and not merely on their ratio. Without getting into lots of algebra, we lose a factor of between  $\sqrt{2}$  and 2 in bandwidth, depending on the details of the frequency compensation scheme, so for a rule of thumb we'll say that

$$
f_{\text{3dB}} \approx \frac{\sqrt{f_{RC} f_T}}{2} \tag{4}
$$

We'll get around 130 kHz transimpedance bandwidth from the LF356 circuit, an improvement of more than 8:1, but still pretty far from 1 MHz.

#### **Noise in the transimpedance amp**

It is obvious from Fig. 3 that all the current sources are treated identically:  $I_d$ ,  $i_{Nshot}$ ,  $i_{Nth}$ , and  $i_{Namp}$  appear in parallel. The Johnson noise  $\overline{i_{Nth}}$  of  $\overline{R}_f$  really appears across  $\overline{R}_f$ , of course, but since the op amp output impedance is low and the currents add linearly, the other end of  $i_{Nth}$  is at ground for noise purposes. As in the simple load resistor case, the rolloff in the frequency response does not degrade the signal-to-current-noise ratio.

The amplifier's voltage noise,  $e_{Namp}$ , is treated differently. Since  $A_1$  is a differential amplifier, we can put  $e_{Namp}$  in either input lead, so we pick the noninverting one because it's easier to analyze. Clearly,  $e_{Namm}$  is multiplied by  $A_1$ 's noninverting gain,

$$
A_{Vcl} = \frac{A_{Vol}}{1 + \frac{A_{Vol}}{1 + j \, V \, C_d Z_f}}
$$
 (5)

where  $Z<sub>f</sub>$  is the complex impedance of the feedback element  $(R_f \text{ in parallel with } C_f)$ . This gain begins to rise at the RC corner frequency of  $C_d$  and  $R_f$ , just where the signal rolloff would have begun if we were using a simple load resistor approach; in fact, the SNR equals that of the same amplifier used as a unity-gain buffer on a photodiode plus load resistor, which is what one would expect. What we've done is tailor the frequency response by using feedback to jiggle the far end of  $R<sub>f</sub>$  but this doesn't get us something for nothing. The addition of  $C_f$  causes  $A_{Vc}$  to level off at  $1/(2\pi R_f C_f)$ .

If  $e_{Namp}$  is very low, or if we are not trying to get a huge bandwidth improvement through the  $(f_T \bullet f_{RC})^{1/2}$  mechanism, this rising noise contribution will not limit us. Otherwise, it will dominate the noise starting at about

$$
f_{3} = \frac{1}{2\pi e_{\text{Namp}} C_{d}} \sqrt{2eI_{d} + i^{2}_{\text{Namp}} + \frac{4kT}{R_{L}}} \tag{6}
$$

Figure 4 shows the noise of our LF356 circuit. It's on linear scales, because loglog noise plots are so deceiving. (Although  $e_{Namp}$  dominates only at the high end, there's a lot more high end than low end.) It only gets worse when we try to go faster this way.

In order that the op amp not dominate the noise, we should choose it by the following rules (worst case specifications apply):

- 1.  $i_{Narm} < 0.5i_{Nth}$ . Make sure the noise of  $R_f$  dominates  $i_{Namp}$ .
- 2.  $e_{\textit{Namp}}$   $< 0.5$   $R_f i_{\textit{Nth}}$ . The same for  $e_{\textit{Namp}}$  in the flatband.
- 3.  $e_{Namp} < 0.5i_{Nth}/(2\pi f_{3\text{ dB}}(C_d + C_{\text{in}})$ . The noise peak should not dominate anywhere in the band.
- 4.  $f_T > 2f_{3dB}^2/f_{RC}$ . The amplifier has to raise the bandwidth enough.



5.  $f_T < 10 f_{3 \text{ dB}}^2/f_{RC}$ . Going too fast risks trouble with ringing and oscillation.

# **No such IC exists: cascode transimpedance amplifiers**

In our case, these rules lead to an amplifier with the following characteristics:  $i_{Namp}$  <  $0.20 \,\, \mathrm{pA/\sqrt{Hz;}}\,\, e_{\mathit{Namp}} < 0.32 \,\, \mathrm{nV/\sqrt{Hz;}}\,\, 250$ MHz  $< f_r < 1250$  MHz. No such amplifier exists.

This is the rock upon which many designs come to grief: the transimpedance amp does nothing whatever to improve the SNR of the photodiode/load resistor combination, it just changes the frequency response. Despite what we may have read, there's nothing inherent or inescapable about this noise peak—it comes from a poor choice of circuit topology that can be fixed.

Another way to reduce the swing across  $C_d$  is to use the common-base transistor amplifier of Fig. 5 (let's just ignore  $R<sub>E</sub>$  for now). Transistor *Q*<sup>1</sup> faithfully transmits its emitter current to its collector, while keeping its emitter at a roughly constant voltage. This idea is called a *cascode*. In the Ebers–Moll transistor model, the small signal resistance  $r_E$  of the transistor's emitter is

$$
r_E = \frac{kT}{eI_C} \tag{7}
$$

where *kT*/*e*is 25 mV at room temperature. Thus, our 2  $\mu$ A photocurrent sees a resistance of 12.5 k $\Omega$ , so that the *RC* bandwidth increases by 8:1 immediately, to about 130 kHz. The summing junction is isolated from  $C_d$ , so that we can raise  $R_f$  to reduce its Johnson noise. On the other hand, we can't improve the bandwidth by using a faster amplifier because the  $(f_{BC} \bullet f_{T})^{1/2}$ mechanism doesn't operate.

Still, we're better off: there are two ways to fix these minor problems while gaining even more bandwidth. And since we know bandwidth isn't everything, let's check the SNR situation.

# **Noise in the cascode**

In the simple load resistor case, the SNR was constant because the signal and all the noise contributions were current sources connected in parallel, so they all rolled off together. Here there is an additional noise contribution from  $Q_1$ , which rises with frequency; it is much more benign than the  $e_{Namp}$  problem with transimpedance amplifiers, however.

A simple noise model of a bipolar junction transistor (BJT) is shown in Fig. 6, which neglects only the Johnson noise of the base resistance  $r_{B}$ <sup>(normally only a</sup> problem when  $I_c \geq 1$  mA). The ideal active device in the model has infinite transconductance and no noise of its own. BJTs operated without feedback exhibit exactly full shot noise in their collector currents, but feedback can suppress this, as we'll see.

Noise current  $i_{nB}$  is the shot noise of the dc base current  $I_B = I_C/\beta_0$ , while  $i_{Nbias}$  is the shot noise of the collector current, which appears in parallel with  $r_F$ . If the emitter is grounded, all of *i<sub>Nbias</sub>* goes from ground into the collector current, and so contributes full shot noise. On the other hand, if the emitter sees a high impedance,  $\emph{\textbf{i}}_{\textit{Nbias}}$ has to flow through  $r_E$ , and none at all winds up in the collector current. The diode's resistance is very large, but the presence of  $C_d$  makes  $i_{Nbias}$  split between  $C_d$ and  $r_F$  by the magnitude ratio of their admittances.

This model gives us the  $Q_1$  contribution to the noise:

$$
i_{NQ_1} = \sqrt{2eI_C} \frac{\omega C_d r_E}{\sqrt{1 + (\omega C_d r_E)^2}} \tag{8}
$$

In an unbiased cascode, where  $R<sub>E</sub>=\infty$  and so  $I_c$  is all from photocurrent, this contribution exactly cancels the *RC* rolloff, giving  $I_c$  exactly full shot noise at all frequencies. Thus, the 1 Hz SNR rolls off exactly as the signal does, and is 3 dB down at the signal corner frequency  $f_c$ —easy to remember, although not a desirable result!

On the other hand, if the applied emitter current  $I_{Eq}$  has only  $\delta$  times full shot noise power, as it will in a minute, the  $i_{Nbias}$ contribution will start to dominate at only

$$
f_{SNR} = f_c \sqrt{\delta} \tag{9}
$$

which turns out to be a serious limitation.

# **Externally biased cascode**

The simpler way of increasing bandwidth is external biasing. Adding a very quiet dc bias current  $I_{Eq}$  to  $I_d$  reduces  $r_E$ , improving  $f_{RC}$ . Choosing  $I_{Eq}$ =20 µA drops  $r_E$  to 1.25  $\overline{k\Omega}$  and increases  $f_{RC}$  to 1.27 MHz. Now the  $C_{in}$  of the op amp becomes the speed limitation. Switching to an LF157 and using  $C_f$ =0.5 pF overcomes  $C_{in}$ , and produces a 1.1 MHz 3 dB bandwidth overall. The collector current now has 10 times less than full shot noise, so (equation 9) predicts that the SNR will be down 3 dB at only 330 kHz, which is not good enough. We could just as easily use  $I_{\text{Eq}}$ =200  $\mu$ A, so that the shot-noise corner would be at 1.3 MHz, but another effect gets in the way, as we'll see.

## **Noise considerations**

Because of the Pauli exclusion principle, currents derived from quiet voltage sources through metal resistors have essentially no shot noise. Resistor  $R<sub>F</sub>$  is connected from the quiet  $V_{bias}$  supply to the slightly-jiggly emitter of  $Q_1$ , so  $I_{Eq}$  will be quiet as long as  $|V_{bias}| >> kT/e$ . Since  $R_E$  has Johnson noise just like  $R_{\vec{t}}$ 's, we also require  $I_dR_F$ >> $kT/e$ , which is a stronger limitation that often requires moderately high supply voltages.

The other important limitation is  $Q_1$ 's base current  $I<sub>b</sub>$ , which has full shot noise. If  $Q_1$ 's dc current gain is  $\beta_0$ , then  $i_{NB}$  limits  $I_{Eq}$  to  $1/\sqrt{\beta_0}$  times full shot noise. You can begin with the near-magical Philips BFG25A/X, but consider using a superbeta transistor ( $\beta \approx 1000$ ) such as an MPSA18.

The calculated transimpedance gain and CNR of the cascoded circuit appear in Fig. 7, with and without an additional 30 -A *IEq*. There's a big improvement in bandwidth and mid-frequency SNR, but the 1 MHz SNR is down by 6 dB due to the bias current noise. Increasing  $I_{Eq}$  makes this problem worse, so we have to look further.

## **Bootstrapping**

When the required value of  $I_{Eq}$  is so large that base current shot noise is a limitation, another technique is superior: bootstrapping. As shown in Fig. 8, driving the cold end of  $D_1$  with a follower  $Q_2$  forces the drop across  $C_d$  to be constant, at least at frequencies where  $X_{C_2}$  is small and  $X_{C_d}$ >>  $I_{E2}$ 

The bootstrap has to have much lower impedance than the cascode, so let's make  $I_{C<sub>2</sub>} >> I_{C<sub>1</sub>}$ . The bootstrap circuit is a bit more complicated to analyze for noise, but the results are nearly the same as for a biased cascode with the same collector current. Assuming  $I_{C2} \gg I_{C1}$ , the noise current from  $Q_2$  flowing to the emitter of  $Q_1$ via  $C_d$  is  $\overline{I_{i}}$   $\overline{I_{i}}$ 

$$
i_{Nbootstrap} = \sqrt{\frac{a}{T_C}} \frac{1}{2} \exp\left[a G_d F_{El} \right]
$$
 (10)

to leading order in  $\omega$ . This is approximately  $(I_{C2}/I_d)^{1/2}$  times smaller than in the unbiased case. It grows linearly with  $\omega$ , so although the bandwidth is increased by  $I_{C2}/I_{d}$ , the SNR is down 3 dB at about  $\omega =$  $(\tilde{I}_{C2}/\tilde{I}_d)^{1/2}/(r_{E2}C_d)$ , just as in the biased cascode case.

Bootstrapping replaces the  $r_{E1}$  of cascode device  $Q_1$  with the  $r_{E2}$  of follower  $Q_2$ , which gives an improvement of  $I_{C_2}/I_{C_1}$ times in bandwidth. By essentially eliminating the capacitive loading on  $Q<sub>1</sub>$ , it also eliminates the effects of  $Q_1$ 's voltage noise.

Bootstrapping suffers voltage-noise multiplication too, but since the *RC* product is not  $R_f C_d$  but  $r_{E_1} C_d$ , a factor of 8 smaller, and the follower's  $e_N$  is usually smaller as well, it is a much less serious problem.

Since current errors are so important, we'll use a superbeta MPSA18 with  $I_{\rm\scriptscriptstyle C2}$ =290  $\mu$ A. The moderately large  $\mathit C_{\scriptscriptstyle e b}$  of this device appears in parallel with  $C_d$ , so it hardly matters; the collector-base capacitance  $C_{ab}$  forms a voltage divider with  $C_{ab}$ , but since it's 50 times smaller, it doesn't matter much either. Altogether, this improves the flatband CNR to 1 dB over shot noise, falling another 2 dB by 1 MHz, and gets us a bandwidth of 2 MHz. The final circuit is shown in Fig. 9, its calculated performance in Fig. 10, and the measured performance of a prototype in Fig. 11, which is somewhat better than the worstcase calculation. The measured shotnoise/dark-noise ratio is 9.5 dB at low frequency, dropping to 4.5 dB at 1 MHz. These numbers correspond to total noise 0.5 dB over shot noise at low frequency, rising to 1 dB over shot noise at 1 MHz.

## **Conclusion**

Now that you've followed all the twists and turns of this article, I hope you're encouraged by the way a couple of inexpensive transistors can sometimes get you a 10:1 bandwidth improvement and lower noise compared with the classical transimpedance amplifier. Next time you're tempted to reach for an expensive APD or analogue-mode PMT, just consider what the right front-end amplifier might do to make your life easier and your product cheaper, more sensitive, and more reliable.

## **References**

1. P. C. D. Hobbs,"Reaching the shot noise limit for \$10," Optics & Photonics News, **2** (4) April, 1991, p. 17.



**Figure 8**. Bootstrapping the unbiased cascode circuit reduces the effects of  $r_F$ , and has performance similar to that of the biased cascode, without the offset current due to R<sub>F</sub>.







**Figure 10**. Performance of the final circuit: CNR is down only 3.3 dB at 1 MHz.



**Figure 11**. Measured performance of the circuit of Fig. 8, showing somewhat better than calculated SNR and bandwidth. Bottom trace: dark noise; top trace: 2 μA *I<sub>d</sub>* added. Measurement setup gain was 2.7.

<sup>2</sup>. See, e.g. Jerald Graeme,"Photodiode Amplifiers: Op Amp Solutions", McGraw-Hill, New York,1995, and Robert A. Pease, "What's all this transimpedance amplifer stuff, anyhow? (Part 1)," Electronic Design, Jan. 8, 2001**.**

**SE Nickols is the CTO of Diablo Mountain Research be reached at SE@DIABLOMOUNTAINRESEARCH.COM**